

Maxwell's Equations for Magnetostatics

From the **point form** of Maxwell's equations, we find that the **static** case reduces to another (in addition to electrostatics) pair of **coupled differential equations** involving magnetic flux density $\mathbf{B}(\bar{r})$ and current density $\mathbf{J}(\bar{r})$:

$$\nabla \cdot \mathbf{B}(\bar{r}) = 0 \quad \nabla \times \mathbf{B}(\bar{r}) = \mu_0 \mathbf{J}(\bar{r})$$

Recall from the **Lorentz force equation** that the magnetic flux density $\mathbf{B}(\bar{r})$ will apply a **force** on current density $\mathbf{J}(\bar{r})$ flowing in volume dV equal to:

$$d\mathbf{F} = (\mathbf{J}(\bar{r}) \times \mathbf{B}(\bar{r})) dV$$

Current density $\mathbf{J}(\bar{r})$ is of course expressed in units of **Amps/meter²**. The units of magnetic flux density $\mathbf{B}(\bar{r})$ are:

$$\frac{\text{Newton} \cdot \text{seconds}}{\text{Coulomb} \cdot \text{meter}} \doteq \frac{\text{Weber}}{\text{meter}^2} \doteq \text{Tesla}$$

- * Recall the units for **electric flux density** $\mathbf{D}(\bar{r})$ are **Colombs/m²**. **Compare** this to the units for **magnetic flux density**—**Webers/m²**.
- * We can say therefore that the units of **electric flux** are **Coulombs**, whereas the units of **magnetic flux** are **Webers**.
- * The concept of **magnetic flux** is much more important and useful than the concept of electric flux, as there is **no** such thing as **magnetic charge**.

We will talk much more later about the concept of **magnetic flux!**

Now, let us consider specifically the **two** magnetostatic equations.

- * First, we note that they specify both the **divergence** and **curl** of magnetic flux density $\mathbf{B}(\bar{r})$, thus **completely** specifying this vector field.
- * Second, it is apparent that the magnetic flux density $\mathbf{B}(\bar{r})$ is **not conservative** (i.e., $\nabla \times \mathbf{B}(\bar{r}) = \mu_0 \mathbf{J}(\bar{r}) \neq 0$).
- * Finally, we note that the magnetic flux density is a **solenoidal** vector field (i.e., $\nabla \cdot \mathbf{B}(\bar{r}) = 0$).

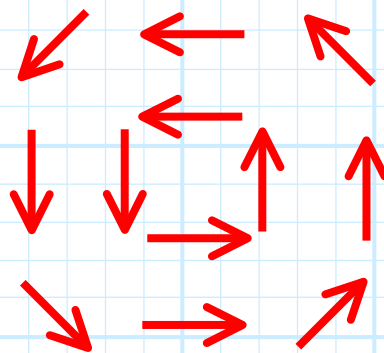
Consider the **first** of the magnetostatic equations:

$$\nabla \cdot \mathbf{B}(\bar{\mathbf{r}}) = 0$$

This equation is sometimes referred to as **Gauss's Law for magnetics**, for its obvious **similarity** to Gauss's Law of electrostatics.

This equation essentially states that the magnetic flux density does **not diverge** nor converge from any point. In other words, it states that there is no such thing as **magnetic charge** !

This of course is **consistent** with our understanding of **solenoidal** vector fields. The vector field will **rotate** about a point, but not diverge from it.



Q: *Just what does the magnetic flux density $\mathbf{B}(\bar{\mathbf{r}})$ rotate around ?*

A: Look at the **second** magnetostatic equation!

The **second** magnetostatic equation is referred to as **Ampere's Circuital Law**:

$$\nabla \times \mathbf{B}(\bar{r}) = \mu_0 \mathbf{J}(\bar{r}) \quad \text{Ampere's Law}$$

This equation indicates that the magnetic flux density $\mathbf{B}(\bar{r})$ **rotates around** current density $\mathbf{J}(\bar{r})$ --the **source** of magnetic flux density is current!.

